

A link between information theory and SPH

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Quick SPH

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$$L = T + \sum_i p_i V_i,$$

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- ▶ some definition of V_i
- ▶ some “closure” to get the pressures p_i

Quick SPH 2

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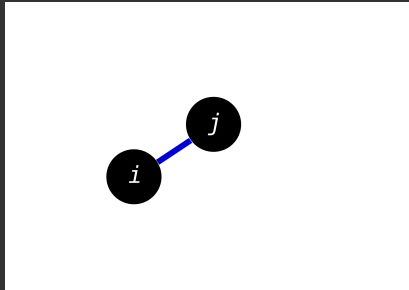
$$V_i = \frac{m_i}{\rho_i} \quad \rho_i = m_j \sum_j W_{ij},$$

$$\mathbf{f}_i = - \sum_j \left(\frac{m_j^2 p_j}{\rho_j^2} + \frac{m_i^2 p_i}{\rho_i^2} \right) \nabla_i W_{ij}.$$

Forget SPH – SFs

All we have is *particles*.

Particle i influences particle j through a **shape function (SF)**,
“ $\phi_i(\mathbf{r}_j)$ ”.

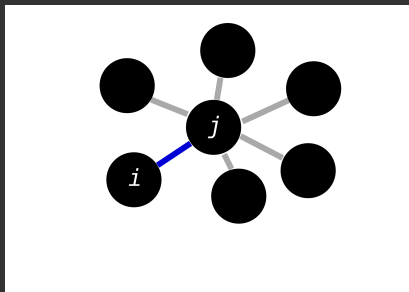


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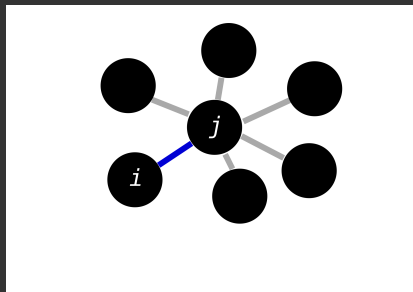
Subject to constraints, such as $\sum_i \phi_i(\mathbf{r}_j) = 1$.



Information entropy

Too many possibilities! But, information theory's “unbiased inference” (Jaynes, 1957) says we should **maximize**

$$S_j = - \sum_i \phi_i(\mathbf{r}_j) \log [\phi_i(\mathbf{r}_j)].$$



SPH?

The resulting SFs are very wide. Hence, Arroyo and Ortiz (2006) added an **energy** to make SFs more local:

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If we forget constraints, the extreme of $F = E - TS$ gives

$$\phi_i(\mathbf{r}_j) = \exp(-\beta r_{ij}^2).$$

More than enough to drive me(*) nuts at the time (*from that time, to be precise.*)

Building a theory: ideal part

A continuum theory already exists[1], so we may inspire ourselves.

$$-S = c^2 \int \rho(\mathbf{r}) \sum_i \phi_i(\mathbf{r}) [\log (\bar{\rho}(\mathbf{r}) \phi_i(\mathbf{r})) - 1] d\mathbf{r}.$$

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Therefore,

$$-S \doteq c^2 \sum_j m_j \sum_i \phi_i(\mathbf{r}_j) \left[\log \left(\frac{\rho_j \phi_i(\mathbf{r}_j)}{m_j} \right) - 1 \right]$$

Building a theory: energy part, constraint

The energy term is

$$E = \beta \int \rho(\mathbf{r}) \sum_i \phi_i(\mathbf{r}) |\mathbf{r} - \mathbf{r}_i|^2 d\mathbf{r}$$

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For a Gaussian kernel, $\beta r_{ij}^2 \propto -\log W_{ij}$, so for any kernel

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Also, a C0 constraint,

$$L_{C0} = \sum_j m_j \alpha_j \left(\sum_i \phi_i(\mathbf{r}_j) - 1 \right)$$

Density, revealed

Before including the pressure: extremezing $F = E - S$ with respect to $\phi_i(\mathbf{r}_j)$ and α_j ,

$$F'' = -c^2 \sum_j m_j \left[\log \left(\frac{m_j Z_j}{\rho_j} \right) + 1 \right]$$

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But, pressure is not included, and we do not want anything happening yet! This **fixes** the density:

$$\rho_j := m_j \sum_i W_{ij}$$

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Finally, there will be a pressure term

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$$V_i = \int \phi_i(\mathbf{r}) d\mathbf{r} \doteq \sum_j \frac{m_j}{\rho_j} \phi_i(\mathbf{r}_j).$$

The C0 model: shape functions

$$\phi_i(r_j) = \frac{W_{ij}z_{i,j}}{Z_j}, \quad z_{i,j} := \exp\left(\frac{p_i}{\rho_j c^2}\right)$$

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$$f_i^{(3)} = - \sum_j \left(\frac{m_j^2 \bar{p}_j}{\rho_j^2} + \frac{m_i^2 \bar{p}_i}{\rho_i^2} \right) \nabla_i W_{ij}, \quad \bar{p}_i := \sum_j \phi_j(r_i) p_j.$$

The C0 model: pressure closure

Since the volumes explicitly depend on the pressures, $V_i = \sum_j \frac{m_j}{\rho_j} \phi_i(\mathbf{r}_j)$, we may look for the pressure field that fixes their values, through a Newton-Raphson method:

$$\sum_j \left(\frac{\partial V_i}{\partial p_j} \right)^{(n)} \left[p_j^{(n+1)} - p_j^{(n)} \right] = - (V_i - V_i^0)^{(n)},$$

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this is a **Poisson equation** to solve for the pressures.

Future: C1 shape functions

$$L_{C1} = \sum_j m_j \lambda_j \cdot \left(\sum_i \phi_i(\mathbf{r}_j) \mathbf{r}_i - \mathbf{r}_j \right)$$

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SFs are now C1 consistent, there is a new force term, but most importantly, perhaps ...

C1 shape functions

The SFs have a sort of boundary detection, and have a “weak-Kronecker condition”.

Thanks, and Ref.

Thanks for your attention.

[1] Daniel Duque, *A unified derivation of Voronoi, power, and finite-element Lagrangian computational fluid dynamics*, European Journal of Mechanics - B/Fluids **98**, 268-278 (2023)

Extra: SFs as interpolants

Consider the smoothing

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Then C0, $\sum_i \phi_i(\mathbf{r}_j) = 1$, means constant fields have proper constant smoothing.

While C1, $\sum_i \phi_i(\mathbf{r}_j) \mathbf{r}_i = \mathbf{r}_j$ means linear fields have proper linear smoothing.